





Session 11: solutions

Exercise 1

- a) The correct variables are $\{n_i\}$ $n_i = 0, \pm 1 \quad \forall i$
with $n = +1 \rightarrow$ full site
 $n = 0 \rightarrow$ empty site

b)
$$H = -\varepsilon \sum_{\langle i,j \rangle} n_i n_j + \mu \sum_i n_i$$

- c) We must map the variables n_i into variables s_i with $s_i = \pm 1$.

Map
$$\begin{aligned} s_i = +1 & \iff n_i = +1 \\ s_i = -1 & \iff n_i = 0 \end{aligned}$$

The relation is

$$n_i = \frac{1}{2} (s_i + 1)$$

Then

$$\begin{aligned} H &= -\varepsilon \sum_{\langle i,j \rangle} \frac{1}{4} (s_i + 1)(s_j + 1) + \mu \sum_i \frac{1}{2} (s_i + 1) = \\ &= -\frac{\varepsilon}{4} \sum_{\langle i,j \rangle} s_i s_j + \left(\frac{\mu}{2} - \frac{\varepsilon z}{4} \right) \sum_i s_i + \text{const} \end{aligned}$$

with z the coordination number of the lattice

Then the mapping with the Ising model is

$$J = \frac{\varepsilon}{4} \quad h = -\frac{\mu}{2} + \frac{\varepsilon z}{4}$$

Exercise 2

Here we go straight to the mapping:

$$s_i = +1 \longrightarrow \text{type A}$$

$$s_i = 0 \longrightarrow \text{vacancy}$$

$$s_i = -1 \longrightarrow \text{type B}$$

of course, other choices are legitimate

then, let's write the chemical potential term:

$$\sum_i \left\{ \frac{\mu_A}{2} s_i (s_i + 1) + \frac{\mu_B}{2} s_i (s_i - 1) \right\} =$$

$= 2$ only if $s_i = +1$, otherwise it is 0

$= 2$ only if $s_i = -1$, otherwise it is 0

$$= \sum_i \left\{ \frac{\mu_A}{2} s_i^2 + \frac{\mu_A}{2} s_i + \frac{\mu_B}{2} s_i^2 - \frac{\mu_B}{2} s_i \right\} =$$

$$= \frac{\mu_A + \mu_B}{2} \sum_i s_i^2 + \frac{\mu_A - \mu_B}{2} \sum_i s_i$$

Now we have learned the trick, we can write the interaction term

$$\sum_{\langle i,j \rangle} \left\{ -\frac{\epsilon_{AA}}{4} s_i(s_i+1) s_j(s_j+1) - \frac{\epsilon_{BB}}{4} s_i(s_i-1) s_j(s_j-1) + \right. \\ \left. - \frac{\epsilon_{AB}}{4} s_i(s_i+1) s_j(s_j-1) \right\} =$$

$$= \sum_{\langle i,j \rangle} \left\{ -\frac{\epsilon_{AA}}{4} s_i^2 s_j^2 - \frac{\epsilon_{AA}}{4} (s_i^2 s_j + s_i s_j^2) - \frac{\epsilon_{AA}}{4} s_i s_j + \right. \\ \left. - \frac{\epsilon_{BB}}{4} s_i^2 s_j^2 + \frac{\epsilon_{BB}}{4} (s_i^2 s_j + s_i s_j^2) - \frac{\epsilon_{BB}}{4} s_i s_j + \right. \\ \left. - \frac{\epsilon_{AB}}{4} s_i^2 s_j^2 + \frac{\epsilon_{AB}}{4} (s_i^2 s_j - s_i s_j^2) + \frac{\epsilon_{AB}}{4} s_i s_j \right\}$$

this is = 0 by exchange symmetry between i and j

Putting everything together:

$$H = - \frac{\epsilon_{AA} + \epsilon_{BB} + \epsilon_{AB}}{4} \sum_{\langle i,j \rangle} s_i^2 s_j^2 - \frac{\epsilon_{AA} - \epsilon_{BB}}{4} \sum_{\langle i,j \rangle} (s_i^2 s_j + s_i s_j^2) + \\ - \frac{\epsilon_{AA} + \epsilon_{BB} - \epsilon_{AB}}{4} \sum_{\langle i,j \rangle} s_i s_j + \frac{\mu_A + \mu_B}{2} \sum_i s_i^2 + \frac{\mu_A - \mu_B}{2} \sum_i s_i$$

This is the most general Hamiltonian with spin 1 and is called the

Blume - Emery - Griffiths model.